# PRODUCTS OF THETA SERIES AND SPECTRAL ANALYSIS 

Dedicated to the memory of Professor Hans Zassenhaus<br>Yuval Z. Flicker and J. G. M. Mars

1. Introduction. The purpose of this note is to propose a new technique in the theory of automorphic forms which will potentially characterize those cusp forms on the general linear group whose symmetric square lifting has a one dimensional constituent. In principle, a cuspidal representation $\pi$ of $\mathbb{G}_{n}=G L(n, \mathbb{A})$, where $\mathbb{A}$ is the ring of adeles of a global field $F$, is parametrized by a complex irreducible representation $\rho$ of dimension $n$ of a form of the Weil group, and the symmetric square lifting $\operatorname{Sym}^{2} \pi$ of $\pi$ is cuspidal precisely when $\operatorname{Sym}^{2} \rho$ is irreducible. A characterization of the $\rho$ such that $\operatorname{Sym}^{2} \rho$ is reducible would suggest a parametrization of the cuspidal $\pi$ whose symmetric square is expected not to be cuspidal, and in particular of the $\pi$ whose symmetric square $L$-function $L\left(s, \pi, S y m^{2}\right)$ - or a twist of it - will not be entire, if $\operatorname{Sym}^{2} \rho$ has a one-dimensional constituent.

An illuminating example is that of a three dimensional $\rho$ with determinant 1 . Its symmetric square is reducible precisely when $\rho$ preserves a quadratic form, and $\rho$ factorizes through the subgroup $(P G L(2, \mathbb{C}) \simeq) S O(3, \mathbb{C})$ of $S L(3, \mathbb{C})$, namely $\rho$ is the symmetric square of some two-dimensional projective representation $\rho_{0}$. This suggests that for a cuspidal representation $\pi$ of $P G L(3, \mathbb{A})$, the $L$-function $L\left(s, \pi, \mathrm{Sym}^{2}\right)$ has a pole precisely when $\pi$ is the symmetric square lifting ([F1], or Gelbart-Jacquet [GJ]) of an automorphic representation of $S L(2, \mathbb{A})$. Patterson and Piatetski-Shapiro [PPS] have shown that the residue of $L\left(s, \varphi, \operatorname{Sym}^{2}\right), \varphi \in \pi$, is $R_{3}(\varphi)=\int_{\mathbb{Z}_{3}^{2} G_{3} \backslash G_{3}} \varphi(g) \Theta(g) \bar{\Theta}(g) d g$ (here $G_{3}=G L(3, F), \mathbb{Z}_{n}=$ center of $\mathbb{G}_{n}$ ), where $\Theta$ are certain "theta" functions on a two-fold covering group of $\mathbb{G}_{3}$. It is then natural to conjecture that the linear form $R$ is non-zero on the cuspidal representation $\pi$ of $\operatorname{PGL}(3, \mathbb{A})$ precisely when it is the symmetric square of a cuspidal representation of $S L(2, \mathbb{A})$. A local analogue of the linear form $R_{3}$ has been studied by Savin $[\mathrm{S}]$ in the unramified case, using the explicit model of the theta representation of [FKS].

Analogous conjectures can be made for all $n$, describing the cuspidal $\pi$ on which $R$ does not vanish (Such $\pi$ might be lifts from $S p_{m}$ if $n=2 m+1$, or $S O_{2 m}$ if $n=2 m$ ). Here we propose a technique to prove these conjectures, by working out the case of $n=2$. This technique is based on applying the theta-kernel to the spectral decomposition of $L^{2}\left(\mathbb{Z}_{2} G_{2} \backslash \mathbb{G}_{2}\right)$. It is likely to generalize to the higher $n$, and in particular to give a new proof of the symmetric square lifting of automorphic forms from $S L(2, \mathbb{A})$ to $\operatorname{PGL}(3, \mathbb{A})$, and a new characterization $\left(R_{3} \neq 0\right.$ on $\pi_{3}=$

[^0]Sym ${ }^{2} \pi_{2}$ ) of the image of the lifting, as well as an extension of the local work of $[\mathrm{S}]$ to the ramified case. But this generalization will require further technical work. We decided to write up the case of $n=2$ to expose our ideas, in the simplest - least technical - case. Our main technical tool, a new type of a summation formula, is described in Proposition 5. Lemmas 2 and 3 deal with the accompanying transfer of orbital integrals.

Consider then a cuspidal representation $\pi$ of $\mathbb{G}=G L(2, \mathbb{A})$. Its symmetric square lifting is an automorphic representation $\operatorname{Sym}^{2} \pi$ of $P G L(3, \mathbb{A})$, whose existence is proven in [GJ] by means of the converse theorem, and in [F1] by means of the trace formula. The $L$-function $L\left(s, \pi, \operatorname{Sym}^{2}\right)=L\left(s, \operatorname{Sym}^{2} \pi\right)$ is entire, but given a character $\chi$ of order two of $\mathbb{A}^{\times} / F^{\times}$, the twisted $L$-function $L\left(s, \pi, \chi \otimes \operatorname{Sym}^{2}\right)=$ $L\left(s, \chi \otimes \operatorname{Sym}^{2} \pi\right)$ will have a pole precisely when $\pi$ is associated with a character $\mu$ of $\mathbb{A}_{E}^{\times} / E^{\times}$, where $E$ is the quadratic separable extension of $F$ defined by $\chi$ using class field theory. It can be shown that the residue of this twisted-by- $\chi L$-function is proportional to $R^{\chi}(\varphi)=\int_{\mathbb{Z}^{2} G \backslash \mathbb{G}} \varphi(g) \Theta(g) \bar{\Theta}^{\chi}(g) d g, \varphi \in \pi$, for suitable $\Theta$-functions on a two-fold covering group of $\mathbb{G}=G L(2, \mathbb{A})$. In fact a similar linear form on $\varphi \in \pi$ appears in [GJ], where $\mathbb{G}$ is replaced by $S L(2, \mathbb{A})$. We use the linear form $R^{\chi}$ to characterize the image of the lifting $\mu \mapsto \pi(\mu)$.

Theorem. Let $\mathbb{A}^{\times}$be the group of ideles of a global field $F$, and $\chi \neq 1$ a quadratic character of $\mathbb{A}^{\times} / F^{\times}$, associated with a quadratic separable field extension $E$ of $F$. Given a character $\mu$ of $\mathbb{A}_{E}^{\times} / E^{\times}$whose restriction to $\mathbb{A}^{\times} / F^{\times}$coincides with $\chi$, there exists a unique automorphic representation $\pi(=\pi(\mu))$ of $P G L(2, \mathbb{A})$, determined as follows. At a place $v$ of $F$ which splits in $E$, there is a character $\mu_{1 v}$ of $F_{v}^{\times}$such that $\mu_{v}((a, b))=\mu_{1 v}(a / b)\left((a, b) \in E_{v}^{\times}=F_{v}^{\times} \times F_{v}^{\times}\right)$. Then the local component $\pi_{v}=$ $\pi\left(\mu_{v}\right)$ of $\pi=\pi(\mu)$ is defined to be the $P G L\left(2, F_{v}\right)$-module $I\left(\mu_{1 v}, \mu_{1 v}^{-1}\right)$ normalizedly induced from the character $\left(\begin{array}{cc}a & * \\ 0 & b\end{array}\right) \mapsto \mu_{1 v}(a / b)$. At a non-split unramified place $v$ of $F$, where $\mu_{v}$ is unramified, there is a character $\mu_{1 v}$ of $F_{v}^{\times}$with $\mu_{v}(z)=\mu_{1 v}(z \bar{z})$. Define $\pi\left(\mu_{v}\right)$ to be $I\left(\mu_{1 v}, \mu_{1 v}^{-1}\right)$. The automorphic representation $\pi(\mu)$ is cuspidal unless $\chi=\eta^{2}$ for some character $\eta$ of $\mathbb{A}^{\times} / F^{\times}$, and $\mu=\bar{\mu}$. In this case $\pi(\mu)$ $=I(\eta, 1 / \eta)$ is a principal series representation. A cuspidal representation $\pi$ of $\operatorname{PGL}(2, \mathbb{A})$ is of the form $\pi(\mu)$ precisely when $R^{\chi}(\varphi)=\int_{\mathbb{Z}^{2} G \backslash \mathbb{G}} \varphi(g) \Theta(g) \bar{\Theta}^{\chi}(g) d g$ is non-zero on $\varphi \in \pi$. In this case, if $\chi$ is a square then $\mu \neq \bar{\mu}$.

The existence of the lifting $\mu \mapsto \pi(\mu)$ is well-known. It was proven using the oscillator representation (see Howe [H], or [MVW]) in Shalika-Tanaka [ST], the converse theorem in Jacquet-Langlands [JL], by stabilizing the trace formula on $S L(2)$ in Labesse-Langlands [LL], by twisting the trace formula by $\chi$ in Kazhdan [K], by quadratic base-change for $G L(2)$ in Langlands [L] (see [F2] for a simpler proof). In all of these works the image of the lifting was characterized by the requirement that $\pi \otimes \chi \simeq \pi$. Our characterization of the image, by the nonvanishing of the form $R^{\chi}$ on $\pi$, is different, and is at the core of our proof. Note
that if exists, $\pi(\mu)$ is uniquely determined by almost all of its components - as specified in the statement of the Theorem - by virtue of the rigidity theorem for $G L(2)$ (see Jacquet-Shalika [JS]).

As noted above, the virtue of the present work is not in proving a new result, or supplying a new proof for an old result. It is in exposing a new method which may extend from the case of $G L(2)$ to the higher rank groups $G L(n)$, $n>2$. In comparison, the method of [ST] - which we proceed to sketch has no known projected extension to $G L(n)$. For simplicity, let us describe the method of [ST] in the case of $S L(2)$. Let $\theta_{1}$ and $\theta_{2}$ be two theta-functions on the two-fold topological central extension $\mathbb{S}$ of $S L(2, \mathbb{A})$. It suffices to show that $(*) \theta_{1}(g) \theta_{2}(g)=\Sigma_{\mu} \phi_{\mu}(g)(g \in S L(2, \mathbb{A}))$, where $\phi_{\mu} \in \pi(\mu)$. Let $\boldsymbol{V}$ be a vector space over $F$ with a quadratic form $q$. Put $\mathbb{V}=\boldsymbol{V}(\mathbb{A})$. Then by $[\mathrm{H}]$ or [MVW], the Schwartz space $C_{c}^{\infty}(\mathbb{V})$ of functions on $\mathbb{V}$ admits commuting representations of $\mathbb{S}$ and the orthogonal group $O(q, \mathbb{V})$ of $q$ on $\mathbb{V}$. If $V=\boldsymbol{V}(F)$ is $F$, and $q(x)=a x^{2}\left(a \in F^{\times}\right)$, one obtains the theta representation of $\mathbb{S}$ on $C_{c}^{\infty}(\mathbb{A})$. If $V=\boldsymbol{V}(F)$ is $E$, and $q(x)=x \bar{x}$ is the norm form on $E$, then one has a direct sum decomposition $C_{c}^{\infty}\left(\mathbb{A}_{E}\right)=\oplus_{\mu} \pi(\mu)$. Since $E=F\left(\tau^{1 / 2}\right) \simeq F \oplus F$ with the quadratic form $q(x, y)=x^{2}-\tau y^{2}$, one has an isomorphism of $S L(2, \mathbb{A})$-modules $C_{c}^{\infty}(\mathbb{A}) \otimes C_{c}^{\infty}(\mathbb{A}) \simeq C_{c}^{\infty}\left(\mathbb{A}_{E}\right)$, and $(*)$ follows (for a complete proof see $[\mathrm{ST}]$, or $[\mathrm{H}]$, [MVW]). To repeat, this method is not known to extend to $G L(n), n>2$.

Our technique might be considered to be conceptually simpler. We consider the well-known spectral and geometric expressions for the kernel of the convolution operator $r(f)$ on $L^{2}(G \backslash \mathbb{G})$ for a Schwartz function $f$ on $\mathbb{G}=P G L(2, \mathbb{A})$, multiply by $\theta_{1}(g) \theta_{2}(g)$, and by a character $\psi(n) \neq 1$ of the upper unipotent subgroup $N \backslash \mathbb{N}$, and integrate over $g \in G \backslash \mathbb{G}$ and over $n \in N \backslash \mathbb{N}$. On the spectral side we get essentially a sum over the cusp forms $(\phi \in) \pi$ of $\mathbb{G}$ of the $R^{\chi}(\pi(f) \phi)$, multiplied by the value at the identity of the Whittaker function of $\phi$. The geometric sum is easily transformed to a sum over $\gamma \in E^{\times}$(rather than $P G L(2, F)!$ ) of the values $f_{E}(\gamma)$ of a function $f_{E}$ in the Schwartz space on $\mathbb{A}_{E}$, transferred from $f$ compatibly with the lifting $\mu \rightarrow \pi(\mu)$ in the unramified case. The Poisson summation formula on $E$ permits writing $\Sigma_{\gamma} f_{E}(\gamma)$ as a sum $\Sigma_{\mu} \mu\left(f_{E}\right)$, and a standard separation argument of "linear independence of characters" establishes the lifting $\mu \rightarrow \pi(\mu)$. This approach extends in principle to $G L(n), n>2$. This we considered interesting, so we thought it was worthwhile to work out carefully the technical details in the test case of $G L(2)$, as a prototype for the general case. This is what we do in this paper.

Let us dispose at once of the degenerate case where there exists a character $\eta$ of $\mathbb{A}^{\times} / F^{\times}$such that $\chi=\eta^{2}$, equivalently $\chi_{v}(-1)=1$ for every place $v$ of $F$, and the character $\mu$ of $\mathbb{A}_{E}^{\times} / E^{\times}$is equal to $\bar{\mu}$, where $\bar{\mu}(x)=\mu(\bar{x}), x \in \mathbb{A}_{E}^{\times}$. Since $\mu=\bar{\mu}$, there is a character $\mu_{1}$ of $\mathbb{A}^{\times} / F^{\times}$such that $\mu=\mu_{1} \circ N$, where $N x=x \bar{x}$ is the norm map from $E$ to $F$. The restriction of $\mu$ to $\mathbb{A}^{\times}$is $\chi$, hence $\mu_{1}^{2}=\chi$. Namely $\chi$ is a square when $\bar{\mu}=\mu$, and we may choose $\eta$ to be $\mu_{1}$. At a place $v$ which splits in $E$, we have $\eta_{v}^{2}=\chi_{v}=1$, hence $\pi_{v}=\pi\left(\mu_{v}\right)$ is by definition the induced $P G L\left(2, F_{v}\right)$-module
$I\left(\eta_{v}, \eta_{v}^{-1}\right)=I\left(\eta_{v}, \eta_{v} \chi_{v}\right)$ (as $\left.\chi_{v}=1\right)$. At a non-split place $v$, by definition $\pi\left(\mu_{v}\right)$ is $I\left(\eta_{v}, \eta_{v}^{-1}\right)=I\left(\eta_{v}, \eta_{v} \chi_{v}\right)$. Hence when $\chi=\eta^{2}$ and $\mu=\bar{\mu}=\eta \circ N$, the character $\mu$ of $\mathbb{A}_{E}^{\times} / E^{\times}$lifts to the principal series (normalizedly induced) representation $I(\eta, \eta \chi)$ of $\operatorname{PGL}(2, \mathbb{A})$.
2. Theta Kernel. Our argument uses the theta-representation of the two-fold cover of the group. For $G L(2)$, an explicit model of this representation is described in $[\mathrm{FM}]$. Let $v$ be a place of $F$, and $\chi_{v}: F_{v}^{\times} \rightarrow \mathbb{C}^{\times}$a unitary character ( $[\mathrm{FM}]$ takes $\chi_{v}=1$, but the general case is similar). Let $C_{\chi_{v}}\left(F_{v}^{\times}\right)$denote the space of smooth functions $u_{v}: F_{v}^{\times} \rightarrow \mathbb{C}$, supported in a compact of $F_{v}$ (if $v$ is finite; having rapid decay at $\infty$ if $v$ is archimedean), which vanish near 0 if $\chi_{v}(-1)=-1$, while if $\chi_{v}(-1)=1$ they have the property that $u_{v 0}(x)=\chi_{v}(t)|t|_{v}^{1 / 2} u_{v}\left(t^{2} x\right)$ is independent of $t$ if $|x|_{v} \leq 1$ and $|t|_{v}$ is sufficiently small (if $v$ is finite; $t \mapsto \chi_{v}(t)|t|_{v}^{1 / 2} u_{v}\left(t^{2} x\right)$ is smooth at $t=0$, and $u_{v 0}(x)$ is defined to be its limit at $t=0$, when $v$ is archimedean). Note that if $\chi_{v}(-1)=1$ then there is a character $\chi_{1 v}$ of $F_{v}^{\times}$with $\chi_{1 v}^{2}=\chi_{v}$, and then $\chi_{1 v} \nu_{v}^{1 / 4} u_{v 0}$ extends to a function on $F_{v}^{\times} / F_{v}^{\times 2}$.

The Weil- or $\theta$-representation of the 2-fold cover $\widetilde{G}_{v}$ of $G_{v}=G L\left(2, F_{v}\right)$ considered in $[\mathrm{FM}]$ acts on $C_{\chi_{v}}\left(F_{v}^{\times}\right)$as follows.

$$
\begin{aligned}
& \left(\theta_{v}\left(s\left(\begin{array}{cc}
z & 0 \\
0 & z
\end{array}\right)\right) u_{v}\right)(\alpha)=(\alpha, z)_{v} \gamma_{v}(z) \chi_{v}(z) u_{v}(\alpha) \quad\left(z, \alpha \in F_{v}^{\times}\right) \\
& \left(\theta_{v}\left(s\left(\begin{array}{ll}
a & 0 \\
0 & 1
\end{array}\right)\right) u_{v}\right)(\alpha)=\chi_{v}(a)|a|_{v}^{1 / 2} u_{v}(a \alpha) \\
& \left(\theta_{v}\left(s\left(\begin{array}{cc}
1 & b \\
0 & 1
\end{array}\right)\right) u_{v}\right)(\alpha)=\psi_{v}\left(\frac{1}{2} b \alpha\right) u_{v}(\alpha) \quad\left(b \in F_{v}\right) \\
& \left(\theta_{v}\left(s\left(\begin{array}{cc}
0 & -1 \\
1 & 0
\end{array}\right)\right) u_{v}\right)(\alpha)=c_{v} \gamma_{v}(\alpha)|\alpha|_{v}^{1 / 2} \int_{F_{v}} \chi_{v}(t)|t|^{1 / 2} u_{v}\left(\alpha t^{2}\right) \psi_{v}(-\alpha t) d t=\left(\mathcal{F} u_{v}\right)(\alpha) .
\end{aligned}
$$

Here $\psi_{v}$ is a non-trivial character of $F_{v}$, and $c_{v}=\gamma_{v}(-1)^{-1 / 2}$ is an eighth root of unity in $\mathbb{C}$. Denote by $R_{v}$ the ring of integers in $F_{v}$, and by $\boldsymbol{\pi}_{v}$ a uniformizer. When $v$ is finite and odd, $\psi_{v}$ has conductor $R_{v}, \chi_{v}$ unramified, $u_{v}^{0}$ is supported on the set of $\varepsilon t^{2}, \varepsilon \in R_{v}^{\times}, t \in R_{v}$, and is given there by $u_{v}^{0}\left(\varepsilon \boldsymbol{\pi}^{2 n}\right)=\chi_{v}(\boldsymbol{\pi})^{n}|\boldsymbol{\pi}|_{v}^{-n / 2}$ $\left(|\varepsilon|_{v}=1, n \geq 0\right)$, then $\theta_{v}(s(k)) u_{v}^{0}=u_{v}^{0}$ for $k \in K_{v}=G L\left(2, R_{v}\right)$.

If $u=\otimes u_{v}$, then [FM] shows that the function

$$
\Theta_{u}^{\chi}(g)=2 \sum_{\alpha \in F^{\times}}(\theta(g) u)(\alpha)+\sum_{\alpha \in F^{\times} / F^{\times 2}}(\theta(g) u)_{0}(\alpha)
$$

on $\widetilde{G}(\mathbb{A})$ is automorphic, namely left-invariant under the discrete subgroup $G=$ $G L(2, F)$ of $\widetilde{G}(\mathbb{A})\left([\mathrm{FM}]\right.$ consider only $\chi=1$; if $\chi_{v}(-1)=-1$ for some $v$ then $\left.(\theta(g) u)_{0} \equiv 0\right)$. Write $\Theta_{u}$ for $\Theta_{u}^{\chi}$ when $\chi=1$. From now on we take a non-trivial character $\chi$ of $\mathbb{A}^{\times} / F^{\times}$of order two, as in the Theorem.

The linear form $R^{\chi}(\varphi)=\int_{\mathbb{Z}^{2} G \backslash \mathbb{G}} \varphi(g) \Theta_{u}(g) \bar{\Theta}_{w}^{\chi}(g) d g\left(u \in C\left(\mathbb{A}^{\times}\right), w \in C_{\chi}\left(\mathbb{A}^{\times}\right)\right)$ appears in the following "spectral" expression on the space $L_{0}^{2}(\mathbb{Z} G \backslash \mathbb{G})$ of cusp forms on $\mathbb{G}=G L(2, \mathbb{A})$ which transform trivially under the center:

$$
\begin{equation*}
\sum_{\substack{\pi \\ \text { cuspidal }}} \sum_{\substack{f \in \operatorname{con} \pi \\ \text { bosimal } \\ \text { basis }}} \int_{\mathbb{Z}^{2} G \backslash \mathbb{G}}(\pi(f) \varphi)(g) \Theta_{u}(g) \bar{\Theta}_{w}^{\chi}(g) d g \cdot \int_{N \backslash \mathbb{N}} \bar{\varphi}(n) \bar{\psi}\left(\frac{1}{2} n\right) d n, \tag{1}
\end{equation*}
$$

where $N=\left\{\left(\begin{array}{ll}1 & * \\ 0 & 1\end{array}\right)\right\}$ and the second integral is the value at $e$ of the Whittaker function $\bar{W}_{\varphi, \psi}$.

We want to show that $\pi$ ranges here over the $\pi=\pi(\mu), \mu: \mathbb{A}_{E}^{\times} / E^{\times} \rightarrow \mathbb{C}^{\times}$, with $\mu \neq \bar{\mu}$ if $\chi$ is a square, and each such $\pi(\mu)$ contributes. For this end note that there are two expressions for the kernel of the convolution operator $r(f)$ on $L^{2}(\mathbb{Z} G \backslash \mathbb{G})$. The geometric expression is $\sum_{\gamma \in Z \backslash G} f\left(g^{-1} \gamma n\right)$. The spectral expression is the sum of the contribution $\sum_{\pi} \sum_{\varphi \in \pi}(\pi(f) \varphi)(g) \bar{\varphi}(n)$ from the cuspidal spectrum, whose integral against $\Theta_{u}(g) \bar{\Theta}_{w}^{\chi}(g) d g \cdot \bar{\psi}\left(\frac{1}{2} n\right) d n$ is (1), and a contribution from the continuous spectrum.
3. Eisenstein Series. The kernel of the operator $r(f)$ on the continuous - nondiscrete - spectrum, takes the form

$$
\frac{1}{\pi} \sum_{\eta} \sum_{\Phi} \int_{i \mathbb{R}} E\left(g, \pi_{s}(f) \Phi, \eta, s\right) \bar{E}(h, \Phi, \eta, s) d s
$$

The first sum ranges over a set of representatives of the classes of characters $\eta$ of $\mathbb{A}^{\times} / F^{\times}$up to multiplication with $\nu^{i s}, s \in \mathbb{R}$, where $\nu(x)=|x|$. The second sum ranges over an orthonormal basis of the space of right smooth functions $\Phi: \mathbb{K} \rightarrow \mathbb{C}$, with $\Phi\left(\left(\begin{array}{cc}a & b \\ 0 & c\end{array}\right) k\right)=\eta(a / c) \Phi(k)$, where $k$ and $\left(\begin{array}{cc}a & b \\ 0 & c\end{array}\right)$ lie in $\mathbb{K}=\Pi K_{v}, K_{v}$ is the standard maximal compact subgroup in $\operatorname{PGL}\left(2, F_{v}\right)$. We trivialize the vector bundle

$$
\begin{aligned}
& I\left(\eta \nu^{s}, \eta^{-1} \nu^{-s}\right)=\left\{\Phi_{s}: \mathbb{G} \rightarrow \mathbb{C} ; \Phi_{s}\left(\left(\begin{array}{cc}
a & b \\
0 & c
\end{array}\right) g\right)\right. \\
& \left.=\eta(a / c)|a / c|^{s+1 / 2} \Phi_{s}(g) ; \quad a, c \in \mathbb{A}^{\times}, b \in \mathbb{A}\right\}
\end{aligned}
$$

via the restriction map $\Phi_{s} \rightarrow \Phi=\Phi_{s} \mid \mathbb{K}$. The Eisenstein series are defined by the sum

$$
E(g, \Phi, \eta, s)=\sum_{\gamma \in B \backslash G} \Phi(\gamma g ; \eta, s),
$$

if $\operatorname{Re}(s)$ is large enough, and by analytic continuation for other $s$ in $\mathbb{C}$. We write $\Phi(g ; \eta, s)$ for $\Phi_{s}(g)$, to emphasize also the dependence on $\eta$.

To compute the integral, $I$, of this kernel against $\Theta_{u}(g) \bar{\Theta}_{w}^{\chi}(g) d g \cdot \bar{\psi}\left(\frac{1}{2} n\right) d n$, we need to recall - and use - the truncation operator $\Lambda^{T}$, where $T>0$ is sufficiently large. If $g=n a k, k \in \mathbb{K}, n \in \mathbb{N}, a=\left(\begin{array}{cc}a & 0 \\ 0 & b\end{array}\right)$, then $d g=|a / b|^{-1}\left(d^{\times} a d^{\times} b\right) d n d k$, and we put $H(g)=|a / b|$. Denote by $\chi(H(g)>T)$ the characteristic function of the $g \in G$ with $H(g)>T$, and similarly with $<$ replacing $>$. The truncation of $\phi$ on $\mathbb{G}$ is

$$
\Lambda^{T} \phi(g)=\phi(g)-\sum_{\delta \in B \backslash G} \phi_{N}(\delta g) \chi(H(\delta g)>T),
$$

where

$$
\phi_{N}(g)=\int_{N \backslash \mathbb{N}} \phi(n g) d n .
$$

The truncation maps slowly increasing to rapidly decreasing functions, and standard arguments imply the following. We have that

$$
I=\frac{1}{\pi} \int_{\mathbb{Z} G \backslash \mathbb{G} N \backslash \mathbb{N}} \int_{\eta} \sum_{\Phi} \sum_{\Phi} \int_{i \mathbb{R}} E\left(g, \pi_{s}(f) \Phi, \eta, s\right) \bar{E}(n, \Phi, \eta, s) d s \Theta_{u}(g) \bar{\Theta}_{w}^{\chi}(g) d g \bar{\psi}\left(\frac{1}{2} n\right) d n
$$

is equal to

$$
I^{\prime}=\frac{1}{\pi} \sum_{\eta} \sum_{\Phi} \lim _{T \rightarrow \infty} \int_{i \mathbb{R}} \int_{\mathbb{Z}^{2} G \backslash \mathbb{G}} \Lambda^{T} E\left(g, \pi_{s}(f) \Phi, \eta, s\right) \Theta_{u}(g) \bar{\Theta}_{w}^{\chi}(g) d g \cdot \bar{E}_{\psi}(\Phi, \eta, s) d s
$$

where

$$
E_{\psi}(\Phi, \eta, s)=\int_{N \backslash \mathbb{N}} E(n, \Phi, \eta, s) \psi\left(\frac{1}{2} n\right) d n
$$

Our aim is to show the following.

1. Lemma. The integral I of the contribution from the continuous spectrum is the sum over the characters $\eta$ of $\mathbb{A}^{\times} / F^{\times}$which satisfy $\eta^{2}=\chi$, of $I_{\eta}$, defined to be the sum over $\Phi$ of

$$
\begin{aligned}
& \frac{1}{4} \int_{\mathbb{K}} \int_{\mathbb{A}^{\times}}\left[F _ { F ^ { \times } } \left[\eta(a)\left(\pi_{0}(f) \Phi(k)+\eta(1 / a)\left(\pi_{0}(f) M \Phi\right)(k)\right] \bar{E}_{\psi}(\Phi, \eta, 0)\right.\right. \\
& \quad \cdot \sum_{\alpha \in F^{\times} / F^{\times 2}}(\theta(k) u)_{0}(a \alpha) \overline{\theta(k) w)_{0}(a \alpha \tau)}|a|^{1 / 2} d^{\times} a d k .
\end{aligned}
$$

Proof. To compute the inner integral, over $g$, in $I^{\prime}$, note that

$$
E_{N}(g, \Phi, \eta, s)=\Phi(g, \eta, s)+(M \Phi)\left(g, \eta^{-1},-s\right)
$$

where $M$ is the standard intertwining operator from $I\left(\eta \nu^{s}, \eta^{-1} \nu^{-s}\right)$ to $I\left(\eta^{-1} \nu^{-s}, \eta \nu^{s}\right)$. Hence

$$
\Lambda^{T} E(g, \Phi, \eta, s)=\sum_{\delta \in B \backslash G}\left[\Phi(\delta g, \eta, s) \chi(H(\delta g)<T)-(M \Phi)\left(\delta g, \eta^{-1},-s\right) \chi(H(\delta g)>T)\right]
$$

and the inner integral, of $\Lambda^{T} E \cdot \Theta_{u} \cdot \bar{\Theta}_{w}^{\chi}$ over $g$, is the difference, which we denote by $J=J(f, \Phi, \eta, s)$,
$\int_{\mathbb{Z} \mathbb{N} B \backslash \mathbb{G}}\left(\pi_{s}(f) \Phi\right)(g, \eta, s) \chi(H(g)<T) A d g-\int_{\mathbb{Z} \mathbb{N} B \backslash \mathbb{G}}\left(M \pi_{s}(f) \Phi\right)\left(g, \eta^{-1},-s\right) \chi(H(g)>T) A d g$.
Here

$$
A=\int_{z \mathbb{Z}^{2} \backslash \mathbb{Z}} \int_{N \backslash \mathbb{N}} \Theta_{u}(n z g) \bar{\Theta}_{w}^{\chi}(n z g) d n d z
$$

Substituting the two sums, over $F^{\times}$and over $F^{\times} / F^{\times 2}$, which define each of $\Theta_{u}$ and $\Theta_{w}^{\chi}$, into $A$, we get the sum of 4 expressions. Note that we may change the order of the summations and the integrations over the compact sets, since for any compact subset $C$ of $\widetilde{G}(\mathbb{A})$, and $u$, there exists a function $u_{1}$ with properties analogous to those of $u$, such that $|(\theta(g) u)(x)| \leq\left|u_{1}(x)\right|$ for all $x \in \mathbb{A}^{\times}$and $g$ in $C$. The same remark applies to the function $w$. In any case, the first term, integrated over the compact $N \backslash \mathbb{N}$, is equal to

$$
\begin{aligned}
& \int_{N \backslash \mathbb{N}} 4 \sum_{\alpha, \beta \in F^{\times}} \psi\left(\frac{1}{2} n(\alpha-\beta)\right)(\theta(z g) u)(\alpha)(\overline{\theta(z g) w})(\beta) d n \\
& =4 \sum_{\alpha \in F^{\times}}(\theta(z g) u)(\alpha)(\overline{\theta(z g) w})(\alpha)=4 \sum_{\alpha \in F^{\times}}(\theta(g) u)(\alpha)(\overline{\theta(g) w})(\alpha) \chi(z) .
\end{aligned}
$$

Integrating this over the compact $Z \mathbb{Z}^{2} \backslash \mathbb{Z}$ we obtain 0 , since $\chi \neq 1$ (is of order two). The second and third terms are similarly shown to be 0 . The fourth term is 0 if there is a place $v$ with $\chi_{v}(-1)=-1$, as then $(\theta(g) w)_{0} \equiv 0$. Suppose then that $\chi_{v}(-1)=1$ for all $v$.

To compute the fourth term in $A$, recall the following action of $\theta$ on $w$.

$$
\begin{aligned}
& \left(\theta\left(s\left(\begin{array}{ll}
1 & b \\
0 & 1
\end{array}\right)\right) w\right)_{0}(\alpha)=w_{0}(\alpha), \quad\left(\theta\left(s\left(\begin{array}{cc}
z & 0 \\
0 & z
\end{array}\right)\right) w\right)_{0}(\alpha)=(\alpha, z) \gamma(z) \chi(z) w_{0}(\alpha), \\
& \left(\theta\left(s\left(\begin{array}{ll}
a & 0 \\
0 & 1
\end{array}\right)\right) w\right)_{0}(\alpha)=\chi(a)|a|^{1 / 2} w_{0}(a \alpha) \\
& \left(\theta\left(s\left(\begin{array}{cc}
a^{2} & 0 \\
0 & 1
\end{array}\right)\right) w\right)_{0}(\alpha)=|a| w_{0}\left(a^{2} \alpha\right)=\chi(a)|a|^{1 / 2} w_{0}(\alpha)
\end{aligned}
$$

since $w_{0}\left(\alpha t^{2}\right)=\chi(t)|t|^{-1 / 2} w_{0}(\alpha)$. The integrand is invariant under $\mathbb{N}$. Its integral over $N \backslash \mathbb{N}$ is its product with $\operatorname{vol}(N \backslash \mathbb{N})=1$. Now, if $\operatorname{ker} \chi=N_{E / F} E^{\times}, E=F(\sqrt{\tau})$,
then $\chi(z)=(\tau, z)$. Indeed, the Hilbert symbol satisfies $\left(-\frac{a}{b}, a+b\right)=(a, b)$, hence $\left(\frac{a^{2}}{\tau b^{2}}, a^{2}-\tau b^{2}\right)=\left(a^{2},-\tau b^{2}\right)=1$, and so $\left(\tau, a^{2}-\theta b^{2}\right)=1$. Carrying out the integration over $z$ in $Z \mathbb{Z}^{2} \backslash \mathbb{Z}$ of the fourth term, we obtain that $A$ is

$$
\begin{aligned}
& \int_{Z \mathbb{Z}^{2} \backslash \mathbb{Z}} \sum_{\alpha, \beta \in F^{\times} / F^{\times 2}}(\alpha, z) \gamma(z)(\beta, z) \bar{\gamma}(z) \chi(z)(\theta(g) u)_{0}(\alpha)(\overline{\theta(g) w})_{0}(\beta) d z \\
= & \sum_{\alpha \in F^{\times} / F^{\times 2}}(\theta(g) u)_{0}(\alpha)\left(\overline{\theta(g) w)_{0}}(\alpha \tau) .\right.
\end{aligned}
$$

We are now in a position to compute the two terms in $J$. Writing $g=n\left(\begin{array}{cc}a t^{2} & 0 \\ 0 & 1\end{array}\right) k$, the first term is

$$
\begin{aligned}
& \int_{\mathbb{K}} \int_{\mathbb{A}^{\times} / F^{\times} \times \mathbb{A}^{\times 2}}\left(\pi_{s}(f) \Phi\right)\left(\left(\begin{array}{ll}
a & 0 \\
0 & 1
\end{array}\right) k, \eta, s\right) \int_{\substack{t \in F^{\times} \backslash \mathbb{A}^{\times} \times \\
\left|t^{2}\right|<T}} \eta\left(t^{2}\right)|t|^{2 s+1} \chi(t)|t|^{1 / 2} \cdot|t|^{1 / 2} \sum_{\alpha \in F^{\times} / F^{\times 2}} \\
& \left(\theta\left(s\left(\begin{array}{ll}
a & 0 \\
0 & 1
\end{array}\right) k\right) u\right)_{0}(\alpha)\left(\theta\left(s\left(\begin{array}{cc}
a & 0 \\
0 & 1
\end{array}\right) k\right) w\right)_{0}(\alpha \tau)|t|^{-2}|a|^{-1} d^{\times} a d^{\times} t d k .
\end{aligned}
$$

The inner integral $\int\left(\eta^{2} \chi\right)(t)|t|^{2 s} d^{\times} t$ over $t \in F^{\times} \backslash \mathbb{A}^{\times},\left|t^{2}\right|<T$, is zero unless $\eta^{2} \chi$ is $\nu^{\lambda}$ for some $\lambda \in i \mathbb{R}$. Replacing $\eta$ by $\eta \nu^{-\lambda / 2}$ we may assume that $\eta^{2}=\chi^{-1}=\chi$, in this case, and then the value of the integral is $\frac{1}{2 s} T^{s}$.

The second term in $J$ is similarly computed, and we conclude that $J$ vanishes unless $\eta^{2}=\chi$, in which case we obtain that $J$ is equal to

$$
\begin{aligned}
& \frac{1}{2 s} T^{s} \int_{\mathbb{K}} \int_{\mathbb{A} \times / F^{\times} \mathbb{A}^{\times 2}}\left(\pi_{s}(f) \Phi\right)\left(\left(\begin{array}{cc}
a & 0 \\
0 & 1
\end{array}\right) k, \eta, s\right) \\
& \quad \sum_{\alpha \in F^{\times} / F^{\times 2}}\left(\theta\left(s\left(\begin{array}{cc}
a & 0 \\
0 & 1
\end{array}\right) k\right) u\right)_{0}(\alpha)\left(\overline{\theta\left(s\left(\begin{array}{cc}
a & 0 \\
0 & 1
\end{array}\right) k\right) w}\right)_{0}(\alpha \tau)|a|^{-s-1} d^{\times} a d k \\
& -\frac{1}{2 s} T^{-s} \int_{\mathbb{K}} \int_{\mathbb{A} \times / F^{\times}}\left(M \pi_{s}(f) \Phi\right)\left(\left(\begin{array}{cc}
a & 0 \\
0 & 1
\end{array}\right) k, \eta^{-1},-s\right) \\
& \quad \sum_{\alpha \in F^{\times} / F^{\times 2}}\left(\theta\left(s\left(\begin{array}{cc}
a & 0 \\
0 & 1
\end{array}\right) k\right) u\right)_{0}(\alpha)\left(\overline{\theta\left(s\left(\begin{array}{cc}
a & 0 \\
0 & 1
\end{array}\right) k\right) w}\right)_{0}(\alpha \tau)|a|^{s-1} d^{\times} a d k .
\end{aligned}
$$

For a given $f$, the sums over $\eta$ and $\Phi$ in $I^{\prime}$ (or $I$ ) are finite. The matrix coefficients $\left(\pi_{s}(f) \Phi, \Phi^{\prime}\right)$ are rapidly decreasing holomorphic functions of $s \in i \mathbb{R}$. Hence $\sum_{\eta} \sum_{\Phi} J(f, \Phi, \eta, s)$ has the form $\frac{1}{2 s} T^{s} h_{1}(s)-\frac{1}{2 s} T^{-s} h_{2}(s)$, where $h_{1}, h_{2}$ are holomorphic on $i \mathbb{R}$, with $h_{1}(0)=h_{2}(0)$. Our lemma now follows from the limit formula

$$
\frac{1}{2 \pi} \lim _{T \rightarrow \infty} \int_{i \mathbb{R}}\left\{\frac{1}{s} T^{s} h_{1}(s)-\frac{1}{s} T^{-s} h_{2}(s)\right\} d s=h_{1}(0)
$$

for such functions $h_{1}, h_{2}$. (The left side is $\lim _{y \rightarrow \infty}(2 \pi i)^{-1} \int_{-\infty}^{\infty}\left[e^{i x y} h_{1}(i x)-e^{-i x y} h_{2}(i x)\right] x^{-1} d x$. Since $\lim _{y \rightarrow \infty} \int_{|x|>1} e^{i x y} h(i x) x^{-1} d x=0$ and $\lim _{y \rightarrow \infty} \int_{|x|<1} e^{i x y}[h(i x)-h(0)] x^{-1} d x=0$, we are left with $(2 \pi i)^{-1} \int_{|x|<1}\left[e^{i x y} h_{1}(0)-e^{-i x y} h_{2}(0)\right] x^{-1} d x=\left(h_{1}(0) / \pi\right) \int_{|x|<1}(\sin x y / x) d x$ $=\left(h_{1}(0) / \pi\right) \underset{|x|<y}{ }(\sin x / x) d x$, which has the limit $h_{1}(0)$ as $y \rightarrow \infty$.)

Remark. It is clear that the integral over $N \backslash \mathbb{N}$ of the kernel of $r(f)$ on the discrete non-cuspidal (one-dimensional) spectrum, multiplied by $\psi\left(\frac{1}{2} n\right)$, is 0 . Hence the one-dimensional automorphic representations do not contribute to our formulae.
4. Geometric Side. We conclude that (1) $+\sum_{\eta^{2}=\chi} I_{\eta}$ is equal to the "geometric sum":

$$
\begin{aligned}
& \int_{\mathbb{Z}^{2} G \backslash \mathbb{G}} \int_{N \backslash \mathbb{N}} \sum_{\gamma \in Z \backslash G} f\left(g^{-1} \gamma n\right) \Theta_{u}(g) \bar{\Theta}_{w}^{\chi}(g) \bar{\psi}\left(\frac{1}{2} n\right) d n d g \\
= & \int_{Z \mathbb{Z}^{2} \mathbb{N} \backslash \mathbb{G}} f_{\psi}\left(g^{-1}\right) \int_{N \backslash \mathbb{N}} \Theta_{u}(n g) \bar{\Theta}_{w}^{\chi}(n g) \bar{\psi}\left(\frac{1}{2} n\right) d n d g,
\end{aligned}
$$

where

$$
f_{\psi}\left(g^{-1}\right)=\int_{\mathbb{N}} f\left(g^{-1} m\right) \bar{\psi}\left(\frac{1}{2} m\right) d m
$$

The inner integral gives

$$
\begin{aligned}
& \int_{N \backslash \mathbb{N}} \Theta_{u}(n g) \bar{\Theta}_{w}^{\chi}(n g) \bar{\psi}\left(\frac{1}{2} n\right) d n=\int_{N \backslash \mathbb{N}} 4 \sum_{\alpha, \beta \in F^{\times}} \psi\left(\frac{1}{2} n(\alpha-\beta-1)\right)(\theta(g) u)(\alpha) \overline{\theta(g) w)}(\beta) d n \\
& \quad+\int_{N \backslash \mathbb{N}} 2 \sum_{\alpha \in F^{\times} / F^{\times 2}} \sum_{\beta \in F^{\times}} \bar{\psi}\left(\frac{1}{2} n(\beta+1)\right)(\theta(g) u)_{0}(\alpha)(\overline{\theta(g) w})(\beta) d n \\
& +\int_{N \backslash \mathbb{N}} 2 \sum_{\beta \in F^{\times} / F^{\times 2}} \sum_{\alpha \in F^{\times}} \psi\left(\frac{1}{2} n(\alpha-1)\right)(\theta(g) u)(\alpha)(\overline{\theta(g) w})_{0}(\beta) d n \\
& =4 \sum_{-1 \neq \beta \in F^{\times}}(\theta(g) u)(\beta+1)(\overline{\theta(g) w})(\beta)+2(\overline{\theta(g) w})(-1) \sum_{\alpha \in F^{\times} / F^{\times 2}}(\theta(g) u)_{0}(\alpha) \\
& \quad+2(\theta(g) u)(1) \sum_{\beta \in F^{\times} / F^{\times 2}}(\overline{\theta(g) w})_{0}(\beta) .
\end{aligned}
$$

Here we used the fact that $\left(\theta\left(\begin{array}{ll}1 & b \\ 0 & 1\end{array}\right) u\right)_{0}(\alpha)=u_{0}(\alpha)$. Consequently, taking
$z \in \mathbb{Z}$, we get

$$
\begin{aligned}
& \int_{N \backslash \mathbb{N}} \Theta_{u}(n z g) \bar{\Theta}_{w}^{\chi}(n z g) \bar{\psi}\left(\frac{1}{2} n\right) d n \\
& =4 \sum_{\beta \in F^{\times}-\{-1\}}(\beta+1, z) \gamma(z)(\theta(g) u)(\beta+1)(\beta, z) \bar{\gamma}(z) \bar{\chi}(z)(\overline{\theta(g) w})(\beta) \\
& \quad+2(-1, z) \bar{\gamma}(z) \bar{\chi}(z)(\overline{\theta(g) w})(-1) \sum_{\alpha \in F^{\times} / F \times 2}(\alpha, z) \gamma(z)(\theta(g) u)_{0}(\alpha) \\
& \quad+2 \gamma(z)(\theta(g) u)(1) \sum_{\beta \in F^{\times} / F^{\times 2}}(\beta, z) \bar{\gamma}(z) \bar{\chi}(z)(\overline{\theta(g) w})_{0}(\beta) .
\end{aligned}
$$

Integrating over $\mathbb{Z} / Z \mathbb{Z}^{2}$ (of course $\left.\mathbb{Z}=\boldsymbol{Z}(\mathbb{A}), Z=\boldsymbol{Z}(F)\right)$ we then get

$$
\begin{aligned}
& (2) \int_{\mathbb{Z}^{2} Z \backslash \mathbb{Z}} \int_{N \backslash \mathbb{N}} \Theta_{u}(n z g) \bar{\Theta}_{w}^{\chi}(n z g) \bar{\psi}\left(\frac{1}{2} n\right) d n d z=4 \sum_{\substack{\beta \in F \times-\{-1\} \\
\frac{\beta+1}{\beta} \in \tau F^{\times 2}}}(\theta(g) u)(\beta+1)(\overline{\theta(g) w})(\beta) \\
& \quad+2(\theta(g) u)_{0}(-\tau)(\overline{\theta(g) w})(-1)+2(\theta(g) u)(1)(\overline{\theta(g) w})_{0}(\tau) .
\end{aligned}
$$

The sum here can be expressed as

$$
\begin{array}{cc}
4 \sum_{\substack{\alpha, \beta \in F^{\times} \\
\alpha \beta \tau \in F^{\times 2}}}(\theta(g) u)(\alpha)(\overline{\theta(g) w})(\beta)=\sum_{\substack{\alpha \in F^{\times} / F^{\times 2} \\
\xi, \eta \in F^{\times} \\
\alpha-\beta=1}}(\theta(g) u)\left(\alpha \xi^{2}\right)(\overline{\theta(g) w})\left(\tau \alpha \eta^{2}\right), \\
& \alpha\left(\xi^{2}-\tau \eta^{2}\right)=1
\end{array}
$$

on replacing $\alpha$ by $\alpha \xi^{2}$ and $\beta$ by $\alpha \tau \eta^{2}$. Assuming that $w=\otimes w_{v} \in C_{\chi}\left(\mathbb{A}^{\times}\right)$, define a function on $\mathbb{A}^{\times} \times \mathbb{A}$ by $w(t, x)=\prod w_{v}\left(t_{v}, x_{v}\right)$, where

$$
w_{v}\left(t_{v}, x_{v}\right)=\chi_{v}\left(x_{v}\right)\left|x_{v}\right|^{1 / 2} w_{v}\left(t_{v} x_{v}^{2}\right) \quad \text { if } x_{v} \neq 0
$$

and

$$
w_{v}\left(t_{v}, 0\right)=\lim _{x_{v} \rightarrow 0} w_{v}\left(t_{v}, x_{v}\right)\left(=w_{v 0}\left(t_{v}\right)\right)
$$

Then $x \mapsto w(t, x)$ is a Schwartz function on $\mathbb{A}$ for every $t$ in $\mathbb{A}^{\times}$, and

$$
w(t, x)=\chi(z)|z|^{1 / 2} w\left(z^{2} t, z^{-1} x\right) \quad\left(t, z \in \mathbb{A}^{\times} ; x \in \mathbb{A}\right)
$$

The analogous definitions - with $\chi=1$ - apply to $u=\otimes u_{v} \in C\left(\mathbb{A}^{\times}\right)$. Then $u(t, x)=|z|^{1 / 2} u\left(z^{2} t, x / z\right)$. If $\alpha, \xi \in F^{\times}$, then $w\left(\alpha \xi^{2}\right)=w(\alpha, \xi)$, and $u\left(\alpha \xi^{2}\right)=$ $u(\alpha, \xi)$. In these notations our sum takes the form

$$
\begin{aligned}
& \sum_{\substack{\alpha \in F^{\times} / F^{\times 2} \\
\xi, \eta \in F^{\times} \\
\alpha\left(\xi^{2}-\tau \eta^{2}\right)=1}}(\theta(g) u)(\alpha, \xi)(\overline{\theta(g) w})(\tau \alpha, \eta) . \\
&
\end{aligned}
$$

Note that when $\eta=0$ we can take $\alpha=1$ (and $\xi= \pm 1$ ), hence the missing term in the last sum is $2(\theta(g) u)(1)(\overline{\theta(g) w})_{0}(\tau)$. When $\xi=0$ we can take $\alpha=-\tau^{-1}$ and $\eta= \pm 1$, hence the corresponding missing term is $2(\theta(g) u)_{0}(-\tau)(\overline{\theta(g) w})(-1)$. These are terms in our integral $\iint \Theta_{u} \bar{\Theta}_{w}^{\chi} \bar{\psi} d n d z$. We conclude that (2) equals

$$
\sum_{\substack{\alpha \in F^{\times} / F^{\times 2} \\ \xi, \eta \in F}}(\theta(g) u)(\alpha, \xi)(\overline{\theta(g) w})(\tau \alpha, \eta)=\sum_{\alpha \in F^{\times} / F^{\times 2}} \sum_{\substack{\gamma \in E^{\times} \\ \alpha N \gamma=1}}(\theta(g) F)(\alpha, \gamma ; u, w)
$$

$$
\alpha\left(\xi^{2}-\tau \eta^{2}\right)=1
$$

where we put $(\theta(g) F)(t, z ; u, w)=\prod_{v}\left(\theta_{v}\left(g_{v}\right) F_{v}\right)\left(t_{v}, z_{v} ; u_{v}, w_{v}\right)$ and

$$
\left(\theta_{v}\left(g_{v}\right) F_{v}\right)\left(t_{v}, z_{v} ; u_{v}, w_{v}\right)=\left(\theta_{v}\left(g_{v}\right) u_{v}\right)\left(t_{v}, x_{v}\right)\left(\overline{\theta_{v}\left(g_{v}\right) w_{v}}\right)\left(\tau t_{v}, y_{v}\right)
$$

$\left(t_{v} \in F_{v}^{\times}, z_{v}=x_{v}+\sqrt{\tau} y_{v} \in E_{v}\right)$. Note that if $\gamma=\xi+\sqrt{\tau} \eta$, then $N \gamma=\xi^{2}-\tau \eta^{2}$. Note also that for $t \in F_{v}, z \in E_{v}$, we have

$$
\begin{aligned}
& \left(\theta_{v}\left(s\left(\begin{array}{cc}
1 & b \\
0 & 1
\end{array}\right)\right) F_{v}\right)(t, z)=\psi_{v}\left(\frac{1}{2} b t N z\right) F_{v}(t, z), \quad b \in F_{v} \\
& \left(\theta_{v}\left(s\left(\begin{array}{cc}
a & 0 \\
0 & 1
\end{array}\right)\right) F_{v}\right)(t, z)=|a|_{v} \chi_{v}(a) F_{v}(a t, z), \quad a \in F_{v}^{\times} \\
& \left(\theta_{v}\left(s\left(\begin{array}{cc}
0 & -1 \\
1 & 0
\end{array}\right)\right) F_{v}\right)(t, z)=\bar{\gamma}_{v}(\tau) \chi_{v}(t)|t|_{v} \int_{E} F_{v}(t, \zeta) \bar{\psi}_{v}\left(\frac{1}{2} t \operatorname{tr}(z \bar{\zeta})\right) d \zeta,
\end{aligned}
$$

and

$$
\theta_{v}\left(s\left(\begin{array}{cc}
a & 0 \\
0 & a
\end{array}\right)\right) F_{v}=F_{v}, \quad a \in F_{v}^{\times} .
$$

Moreover

$$
F(t, z)=|s| \chi(s) F\left(s^{2} t, z / s\right) \quad\left(s, t \in \mathbb{A}^{\times} ; z \in \mathbb{A}_{E}\right)
$$

Note that $\theta_{v}$ is a representation of the group $G L\left(2, F_{v}\right)$ itself (even of $P G L\left(2, F_{v}\right)$ ) on the space of the functions $F_{v}\left(t_{v}, z_{v}\right)$, which are smooth on $F_{v}^{\times} \times E_{v}$. Hence (2) can be written as

$$
2 \sum_{\gamma \in E^{\times} / F^{\times}}(\theta(g) F)\left(N \gamma^{-1}, \gamma ; u, w\right)
$$

and the total "geometric sum" is equal to

$$
2 \sum_{\gamma \in E^{\times} / F^{\times}} \int_{\mathbb{Z} \mathbb{N} \backslash \mathbb{G}} f_{\psi}\left(g^{-1}\right)(\theta(g) F)\left(N \gamma^{-1}, \gamma ; u, w\right) d g
$$

For $z \in E_{v}^{\times}$, define

$$
f_{E_{v}}(z)=|z \bar{z}|_{F_{v}}^{-1 / 2} \int_{Z_{v} N_{v} \backslash G_{v}} f_{v, \psi}\left(g^{-1}\right)\left(\theta_{v}(g) F_{v}\right)\left(N z^{-1}, z ; u_{v}, w_{v}\right) d g
$$

The function $f_{E_{v}}$ on $E_{v}^{\times}$satisfies $f_{E_{v}}(a z)=\chi_{v}(a) f_{E_{v}}(z)\left(a \in F_{v}^{\times}, z \in E_{v}^{\times}\right)$, and $f_{E_{v}}(\bar{z})=\chi_{v}(-1) f_{E_{v}}(z)$.

## 5. Transfer of Functions.

2. Lemma. The function $f_{E_{v}}(\gamma)$ extends to a smooth function on $E_{v}^{\times}$.

Proof. This is clear, since the function $F_{v}$ is smooth on $F_{v}^{\times} \times E_{v}$, and the integration ranges over a compact set of $g$, depending on $f_{v}$. Alternatively stated, let us pass to local notations - drop $v$ - to simplify the notations. Write $\gamma=x+\sqrt{\tau} y$. Then up to a factor which is smooth in $\gamma$, our expression is the integral over $k$ in $K$ of

$$
\chi(y)|x y|^{1 / 2} \int_{F^{\times}} \chi(a) f_{\psi}\left(k^{-1}\left(\begin{array}{cc}
a^{-1} & 0 \\
0 & 1
\end{array}\right)\right)(\theta(k) u)\left(\frac{a x^{2}}{\gamma \bar{\gamma}}\right)(\overline{\theta(k) w})\left(\frac{\tau a y^{2}}{\gamma \bar{\gamma}}\right) d^{\times} a
$$

at $\gamma$ with $x y \neq 0$. The only possible points where this may not be smooth are at $x=0$ or $y=0$. But at these points we have that

$$
(\theta(k) u)\left(\frac{a x^{2}}{\gamma \bar{\gamma}}\right)=|x|^{-1 / 2}(\theta(k) u)_{0}(a / \gamma \bar{\gamma}) \quad(|x| \text { small })
$$

and

$$
(\theta(k) w)\left(\frac{\tau a y^{2}}{\gamma \bar{\gamma}}\right)=\chi(y)|y|^{-1 / 2}(\theta(k) w)_{0}(\tau a / \gamma \bar{\gamma}) \quad(|y| \text { small })
$$

hence the lemma again follows.
Consider next the case of a spherical function, first at a place $v$ which splits in $E$.
3. Lemma. If $v$ splits in $E, \psi_{v}$ has conductor $R_{v}, u_{v}=u_{v}^{0}, w_{v}=w_{v}^{0}$, and $f_{v}$ is spherical, then $f_{E_{v}}((a, b))=F_{f_{v}}\left(\left(\begin{array}{cc}a & 0 \\ 0 & b\end{array}\right)\right) ; a, b \in F_{v}^{\times}$.

Corollary. A character $\mu_{v}$ of $E_{v}^{\times} / F_{v}^{\times}$is of the form $\mu_{v}((a, b))=\mu_{1 v}(a / b)$ for some character $\mu_{1 v}$ of $F_{v}^{\times}$, and we have $\mu_{v}\left(f_{E_{v}}\right)=\operatorname{tr} I\left(\mu_{1 v}, \mu_{1 v}^{-1} ; f_{v}\right)$ for every spherical function $f_{v}$ on $\operatorname{PGL}\left(2, F_{v}\right)$.

Proof. Note that

$$
\mu_{v}\left(f_{E_{v}}\right)=\int_{E_{v}^{\times} / F_{v}^{\times}} \mu_{v}((a, b)) f_{E_{v}}((a, b)) d^{\times}(a / b)
$$

is equal to

$$
\int_{F_{v}^{\times}} \mu_{1 v}(a / b) F_{f_{v}}\left(\left(\begin{array}{cc}
a & 0 \\
0 & b
\end{array}\right)\right) d^{\times}(a / b)=\operatorname{tr} I\left(\mu_{1 v}, \mu_{1 v}^{-1} ; f_{v}\right) .
$$

Proof of Lemma. Suppose that $v$ splits in $E$, thus $E_{v}=F_{v} \oplus F_{v}$ and $\bar{\gamma}=(d, c)$ if $\gamma=$ $(c, d)$, and assume that $\psi_{v}$ has conductor $R_{v}, u_{v}=u_{v}^{0}, w_{v}=w_{v}^{0}$, and $f_{v}$ is spherical.

Note that $\chi_{v}=1$. Then using the Iwasawa decomposition $d g=d n d^{\times} a /|a| d k$, and noting that $\theta(k) u_{v}^{0}=u_{v}^{0}$, and $\theta(k) w_{v}^{0}=w_{v}^{0}$, we obtain the following expression for $f_{E_{v}}((c, d))$ :

$$
|c d|_{v}^{-1 / 2} \int_{F_{v}^{\times}} f_{v, \psi}\left(\left(\begin{array}{cc}
a^{-1} & 0 \\
0 & 1
\end{array}\right)\right) u_{v}^{0}\left(\left(a / c d, \frac{1}{2}(c+d)\right)\right) \bar{w}_{v}^{0}((\tau a / c d,(c-d) / 2 \sqrt{\tau})) d^{\times} a .
$$

Here $\gamma=(c, d)$ in $E_{v}^{\times}$can be expressed as $x_{v}+\sqrt{\tau} y_{v}$, where $\sqrt{\tau}=(\sqrt{\tau},-\sqrt{\tau})$ and $x_{v}=\frac{1}{2}(\gamma+\bar{\gamma})=\frac{1}{2}(c+d)$, and $y_{v}=(\gamma-\bar{\gamma}) / 2 \sqrt{\tau}=(c-d) / 2 \sqrt{\tau}$. At $\gamma \neq \pm \bar{\gamma}$ in $E_{v}^{\times}$, we obtain

$$
\left|\frac{c^{2}-d^{2}}{c d}\right|_{v}^{1 / 2} \int_{F_{v}^{\times}} f_{v, \psi}\left(\left(\begin{array}{cc}
a^{-1} & 0 \\
0 & 1
\end{array}\right)\right) u_{v}^{0}\left(\frac{a(c+d)^{2}}{4 c d}\right) \bar{w}_{v}^{0}\left(\frac{a(c-d)^{2}}{4 c d}\right) d^{\times} a .
$$

This expression is not changed if $(c, d)$ is replaced by $(d, c)$ or $(-c, d)$, and $(c, d)$ is taken modulo $F_{v}^{\times}$. We may take then $d=1$ and $|c| \leq 1$. Consider first the case that $|c|=1, c \neq \pm 1$. We may assume that $|c+1|=1$, and $|c-1| \leq 1$. Then our expression is

$$
\begin{aligned}
& |c-1|_{v}^{1 / 2} \int_{F_{v}^{\times}} f_{v, \psi}\left(\left(\begin{array}{cc}
a^{-1} & 0 \\
0 & 1
\end{array}\right)\right) u_{v}^{0}(a) \bar{w}_{v}^{0}\left(a(c-1)^{2}\right) d^{\times} a \\
& =\int_{F_{v}^{\times}} f_{v, \psi}\left(\left(\begin{array}{cc}
a^{-1} & 0 \\
0 & 1
\end{array}\right)\right) u_{v}^{0}(a) \bar{w}_{v}^{0}(a) d^{\times} a .
\end{aligned}
$$

If $|c|=\left|\boldsymbol{\pi}^{n}\right|<1$, our expression is

$$
\left|\boldsymbol{\pi}^{n}\right|^{-1 / 2} \int_{F_{v}^{\times}} f_{v, \psi}\left(\left(\begin{array}{cc}
a^{-1} & 0 \\
0 & 1
\end{array}\right)\right) u_{v}^{0}\left(a \boldsymbol{\pi}^{-n}\right) \bar{w}_{v}^{0}\left(a \boldsymbol{\pi}^{-n}\right) d^{\times} a
$$

this last expression is then valid for $n=0$ too. The integrand is non-zero only when $a \in \pi^{n+2 m} R_{v}^{\times}, m \geq 0$, and we get

$$
\begin{aligned}
& \left|\boldsymbol{\pi}^{n}\right|^{-1 / 2} \sum_{m \geq 0} f_{v, \psi}\left(\left(\begin{array}{cc}
\boldsymbol{\pi}^{-n-2 m} & 0 \\
0 & 1
\end{array}\right)\right)|\boldsymbol{\pi}|^{-m} \\
& =\sum_{m \geq 0} q^{m+\frac{1}{2} n}\left[f_{v}\left(\left(\begin{array}{cc}
\pi^{-n-2 m} & 0 \\
0 & 1
\end{array}\right)\right)-f_{v}\left(\left(\begin{array}{cc}
\boldsymbol{\pi}^{-n-2 m-2} & 0 \\
0 & 1
\end{array}\right)\right)\right] \\
& =q^{\frac{1}{2} n}\left[f_{v}\left(\left(\begin{array}{cc}
\boldsymbol{\pi}^{-n} & 0 \\
0 & 1
\end{array}\right)\right)+\sum_{m \geq 1}\left(1-\frac{1}{q}\right) q^{m} f_{v}\left(\left(\begin{array}{cc}
\pi^{-n-2 m} & 0 \\
0 & 1
\end{array}\right)\right)\right] .
\end{aligned}
$$

Note that

$$
\begin{aligned}
& F_{f_{v}}\left(\left(\begin{array}{ll}
a & 0 \\
0 & b
\end{array}\right)\right)=\frac{|a-b|_{v}}{|a b|_{v}^{1 / 2}} \int_{F_{v}} f_{v}\left(\left(\begin{array}{cc}
1 & -x \\
0 & 1
\end{array}\right)\left(\begin{array}{cc}
a & 0 \\
0 & b
\end{array}\right)\left(\begin{array}{ll}
1 & x \\
0 & 1
\end{array}\right)\right) d x \\
& =|a / b|_{v}^{1 / 2} \int_{F_{v}} f_{v}\left(\left(\begin{array}{cc}
a & 0 \\
0 & b
\end{array}\right)\left(\begin{array}{cc}
1 & x \\
0 & 1
\end{array}\right)\right) d x
\end{aligned}
$$

and so, for our spherical $f_{v}$, we have

$$
\begin{aligned}
& F_{f_{v}}\left(\left(\begin{array}{cc}
\boldsymbol{\pi}^{-n} & 0 \\
0 & 1
\end{array}\right)\right)=\left|\boldsymbol{\pi}^{-n}\right|_{v}^{1 / 2} \int f_{v}\left(\left(\begin{array}{cc}
\boldsymbol{\pi}^{-n} & 0 \\
0 & 1
\end{array}\right)\left(\begin{array}{ll}
1 & x \\
0 & 1
\end{array}\right)\right) d x \\
& =\left|\boldsymbol{\pi}^{-n}\right|^{1 / 2}\left[f_{v}\left(\left(\begin{array}{cc}
\boldsymbol{\pi}^{-n} & 0 \\
0 & 1
\end{array}\right)\right)+\int_{|x|>1} f_{v}\left(\left(\begin{array}{cc}
\boldsymbol{\pi}^{-n} x^{2} & 0 \\
0 & 1
\end{array}\right)\right) d x\right] \\
& =q^{n / 2}\left[f_{v}\left(\left(\begin{array}{cc}
\boldsymbol{\pi}^{-n} & 0 \\
0 & 1
\end{array}\right)\right)+\left(1-\frac{1}{q}\right) \sum_{m \geq 1} q^{m} f_{v}\left(\left(\begin{array}{cc}
\pi^{-n-2 m} & 0 \\
0 & 1
\end{array}\right)\right)\right]
\end{aligned}
$$

Here we used the decomposition

$$
\begin{aligned}
& \left(\begin{array}{cc}
1 & 0 \\
0 & \pi^{-n}
\end{array}\right)\left(\begin{array}{ll}
1 & 0 \\
x & 1
\end{array}\right)=\left(\begin{array}{cc}
1 & 0 \\
0 & \pi^{-n}
\end{array}\right)\left(\begin{array}{cc}
1 & \frac{1}{x} \\
0 & 1
\end{array}\right)\left(\begin{array}{cc}
x^{-1} & 0 \\
0 & x
\end{array}\right)\left(\begin{array}{cc}
0 & -1 \\
1 & 0
\end{array}\right)\left(\begin{array}{cc}
1 & 1 / x \\
0 & 1
\end{array}\right) \\
& =\left(\begin{array}{cc}
1 & \boldsymbol{\pi}^{n} / x \\
0 & 1
\end{array}\right)\left(\begin{array}{cc}
x^{-1} & 0 \\
0 & x \boldsymbol{\pi}^{-n}
\end{array}\right)\left(\begin{array}{cc}
0 & -1 \\
1 & 0
\end{array}\right)\left(\begin{array}{cc}
1 & 1 / x \\
0 & 1
\end{array}\right) \equiv\left(\begin{array}{cc}
x^{-1} & 0 \\
0 & x \boldsymbol{\pi}^{-n}
\end{array}\right) \quad \text { if }|x| \geq 1
\end{aligned}
$$

The lemma follows.
6. Non-split Case. Suppose now that $E_{v} / F_{v}$ is an unramified field extension, the conductor of $\psi_{v}$ is $R_{v}$ (thus $\psi_{v}=1$ on $R_{v}$ but $\left.\psi_{v}\left(\boldsymbol{\pi}_{v}^{-1} R_{v}\right) \neq 1\right),|2|_{v}=1, f_{v}$ is spherical ( $K_{v}=G L\left(2, R_{v}\right)$-biinvariant), and $u_{v}=u_{v}^{0}, w_{v}=w_{v}^{0}$, the $K_{v}$-invariant elements in $C\left(F_{v}^{\times}\right)$and $C_{\chi_{v}}\left(F_{v}^{\times}\right)$. Then $d g=d n \cdot|a|^{-1} d^{\times} a \cdot d k$ if $g=n\left(\begin{array}{cc}a & 0 \\ 0 & 1\end{array}\right) k$ in $Z_{v} \backslash G_{v}=N_{v} A_{v} K_{v}$. Hence

$$
\begin{aligned}
& f_{E_{v}}(\gamma)=|\gamma \bar{\gamma}|_{v}^{-1 / 2} \int_{F_{v}^{\times}} f_{v, \psi}\left(\left(\begin{array}{cc}
a^{-1} & 0 \\
0 & 1
\end{array}\right)\right) \\
& \cdot \chi_{v}(a) u_{v}^{0}\left(a / \gamma \bar{\gamma}, \frac{1}{2}(\gamma+\bar{\gamma})\right) \bar{w}_{v}^{0}\left(\tau a / \gamma \bar{\gamma}, \frac{\gamma-\bar{\gamma}}{2 \sqrt{\tau}}\right) d^{\times} a .
\end{aligned}
$$

This is a function on $E_{v}^{\times}$which transforms under $F_{v}^{\times}$according to $\chi_{v}$. Hence we may assume that $\gamma \bar{\gamma}$ and $\tau$ are units in $F_{v}^{\times}$. At $\gamma$ with $\bar{\gamma} \neq \pm \gamma$, we put $x=(\gamma+\bar{\gamma}) / 2$, $y=(\gamma-\bar{\gamma}) / 2 \sqrt{\tau}$, and then

$$
\begin{aligned}
f_{E_{v}}(\gamma) & =|x y|_{F_{v}}^{1 / 2} \chi_{v}(y) \int_{F_{v}^{\times}} f_{v, \psi}\left(\left(\begin{array}{cc}
a^{-1} & 0 \\
0 & 1
\end{array}\right)\right) \chi_{v}(a) u_{v}^{0}\left(a x^{2}\right) \bar{w}_{v}^{0}\left(a y^{2}\right) d^{\times} a \\
& =\int_{F_{v}^{\times}} f_{v, \psi}\left(\left(\begin{array}{cc}
a^{-1} & 0 \\
0 & 1
\end{array}\right)\right) u_{v}^{0}(a) \bar{w}_{v}^{0}(a) d^{\times} a .
\end{aligned}
$$

The last integral ranges over $R_{v}^{\times} \pi^{2 m}, m \geq 0$, where $\chi_{v}(a)=1$, and we used $u_{v}^{0}\left(a x^{2}\right)=u_{v}^{0}(a)|x|_{v}^{-1 / 2}$ and $w_{v}^{0}\left(a y^{2}\right)=|y|_{v}^{-1 / 2} \chi_{v}(y) w_{v}^{0}(a)$ for the last equality. It follows that in the unramified-spherical case, $f_{E_{v}}(\gamma)$ depends only on the parity of
the valuation of $\gamma$. If - moreover $-f_{v}$ is the unit element $f_{v}^{0}$ of the Hecke algebra, and $|a|_{v} \leq 1$, then

$$
\begin{aligned}
& f_{v, \psi}^{0}\left(\left(\begin{array}{cc}
a^{-1} & 0 \\
0 & 1
\end{array}\right)\right)=\int_{F_{v}} f_{v}^{0}\left(\left(\begin{array}{cc}
a^{-1} & 0 \\
0 & 1
\end{array}\right)\left(\begin{array}{ll}
1 & x \\
0 & 1
\end{array}\right)\right) \bar{\psi}\left(\frac{1}{2} x\right) d x \\
& =f_{v}^{0}\left(\left(\begin{array}{cc}
a^{-1} & 0 \\
0 & 1
\end{array}\right)\right)-f_{v}^{0}\left(\left(\begin{array}{cc}
a^{-1} \pi^{-2} & 0 \\
0 & 1
\end{array}\right)\right)
\end{aligned}
$$

is equal to $f_{v}^{0}(I)=1$ if $|a|_{v}=1$. Hence $f_{E_{v}}^{0}(\gamma) \equiv(-1)^{v^{\text {al }}{ }_{E_{v}}(\gamma)}$.
4. Lemma. When $E_{v} / F_{v}$ is an unramified field extension, $f_{v}$ is spherical, $\psi_{v}$ has conductor $R_{v},|2|_{v}=1, u_{v}=u_{v}^{0}, w_{v}=w_{v}^{0}$, and $\mu_{1 v}$ is the unramified character of $F_{v}^{\times}$whose value at the uniformizer $\boldsymbol{\pi}_{v}$ of $R_{v}$ is $i=\sqrt{-1}$, then

$$
\mu_{v}\left(f_{E_{v}}\right)=\operatorname{trI}\left(\mu_{1 v}, \mu_{1 v}^{-1} ; f_{v}\right)
$$

where $\mu_{v}$ denotes the unramified "sign" character of $E_{v}^{\times}$, whose value at a uniformizer $\boldsymbol{\pi}_{v}$ of $R_{E_{v}}$ is -1 . Here $\mu_{v}\left(f_{E_{v}}\right)=\int_{E_{v}^{\times} / F_{v}^{\times}} \mu_{v}(\gamma) f_{E_{v}}(\gamma) d^{\times} \gamma$ is the value of $f_{E_{v}}$ at a $\gamma$ in $R_{E_{v}}^{\times}$.

Proof. We drop the index $v$ to simplify the notations, and recall that $f_{E}(\gamma), \gamma \in R_{E}^{\times}$, is independent of $\gamma$, and is given by

$$
\begin{aligned}
f_{E}(\gamma) & =\int_{F \times}\left[f\left(\left(\begin{array}{cc}
a^{-1} & 0 \\
0 & 1
\end{array}\right)\right)-f\left(\left(\begin{array}{cc}
a^{-1} \boldsymbol{\pi}^{-2} & 0 \\
0 & 1
\end{array}\right)\right)\right] u^{0}(a) \bar{w}^{0}(a) d^{\times} a \\
& =\sum_{n \geq 0}\left[f\left(\left(\begin{array}{cc}
\boldsymbol{\pi}^{-2 n} & 0 \\
0 & 1
\end{array}\right)\right)-f\left(\left(\begin{array}{cc}
\boldsymbol{\pi}^{-2 n-2} & 0 \\
0 & 1 n
\end{array}\right)\right)\right]|\boldsymbol{\pi}|^{-n} \chi\left(\boldsymbol{\pi}^{n}\right) \\
& =f(I)+\left(1+\frac{1}{q}\right) \sum_{n \geq 1} f\left(\left(\begin{array}{cc}
\boldsymbol{\pi}^{-2 n} & 0 \\
0 & 1
\end{array}\right)\right)(-q)^{n}, \quad q=|\boldsymbol{\pi}|^{-1} .
\end{aligned}
$$

On the other hand, since $\mu_{1}(\boldsymbol{\pi})=i$, we have

$$
\begin{aligned}
& \operatorname{tr} I\left(\mu_{1}, \mu_{1}^{-1} ; f\right)=\sum_{n \in \mathbb{Z}} F_{f}\left(\left(\begin{array}{cc}
\pi^{-n} & 0 \\
0 & 1
\end{array}\right)\right) i^{n}=F_{f}(I)+\sum_{n \geq 1} F_{f}\left(\left(\begin{array}{cc}
\pi^{-n} & 0 \\
0 & 1
\end{array}\right)\right)\left(i^{n}+i^{-n}\right) \\
& =F_{f}(I)+2 \sum_{n \geq 1}(-1)^{n} F_{f}\left(\left(\begin{array}{cc}
\pi^{-2 n} & 0 \\
0 & 1
\end{array}\right)\right)
\end{aligned}
$$

But

$$
F_{f}\left(\left(\begin{array}{cc}
\boldsymbol{\pi}^{-n} & 0 \\
0 & 1
\end{array}\right)\right)=q^{n / 2}\left[f\left(\left(\begin{array}{cc}
\boldsymbol{\pi}^{-n} & 0 \\
0 & 1
\end{array}\right)\right)+\left(1-\frac{1}{q}\right) \sum_{m \geq 1} q^{m} f\left(\left(\begin{array}{cc}
\pi^{-n-2 m} & 0 \\
0 & 1
\end{array}\right)\right)\right]
$$

So we get

$$
\begin{aligned}
& =F_{f}(I)+2 \sum_{n \geq 1}(-1)^{n} F_{f}\left(\left(\begin{array}{cc}
\pi^{-2 n} & 0 \\
0 & 1
\end{array}\right)\right)=f(I)+\left(1-\frac{1}{q}\right) \sum_{m \geq 1} q^{m} f\left(\left(\begin{array}{cc}
\boldsymbol{\pi}^{-2 m} & 0 \\
0 & 1
\end{array}\right)\right) \\
& +2 \sum_{k \geq 1}(-q)^{k}\left[f\left(\left(\begin{array}{cc}
\pi^{-2 k} & 0 \\
0 & 1
\end{array}\right)\right)+\left(1-\frac{1}{q}\right) \sum_{m \geq 1} q^{m} f\left(\left(\begin{array}{cc}
\pi^{-2 k-2 m} & 0 \\
0 & 1
\end{array}\right)\right)\right] \\
& =f(I)+\sum_{n \geq 1} f\left(\left(\begin{array}{cc}
\pi^{-2 n} & 0 \\
0 & 1
\end{array}\right)\right)\left[\left(1-\frac{1}{q}\right) q^{n}+2(-q)^{n}+2\left(1-\frac{1}{q}\right) q^{n} \sum_{1 \leq k<n}(-1)^{k}\right] \\
& =f(I)+\sum_{n \geq 1}\left(1+\frac{1}{q}\right)(-q)^{n} f\left(\left(\begin{array}{cc}
\pi^{-2 n} & 0 \\
0 & 1
\end{array}\right)\right)=f_{E}(\gamma),
\end{aligned}
$$

as asserted.
7. Conclusion. So far we have shown the following.
5. Proposition. Given a finite set $V$ of places of $F$ containing the archimedean places and those which ramify in $E$, and those where the conductor of $\psi$ is not $R_{v}$, and those where $u_{v} \neq u_{v}^{0}$, or $w_{v} \neq w_{v}^{0}$, for any test function $f=\otimes f_{v}, f_{v} \in$ $C_{c}^{\infty}\left(Z_{v} \backslash G_{v}\right), f_{v}$ is spherical $\left(K_{v}=G L\left(2, R_{v}\right)\right.$-biinvariant) for all $v \notin V$, and $f_{v}=f_{v}^{0}\left(=\right.$ characteristic function of $\left.Z_{v} K_{v}\right)$ for almost all $v$, we have the equality

$$
(1)+\sum_{\eta^{2}=\chi} I_{\eta}=2 \sum_{\gamma \in E^{\times} / F^{\times}} f_{E}(\gamma)=\sum_{\mu} \mu\left(f_{E}\right) .
$$

Here $I_{\eta}$ is defined in Lemma 1. Moreover, $f_{E}(a)=\prod_{v} f_{E_{v}}\left(a_{v}\right)$ for $a=\left(a_{v}\right) \in \mathbb{A}_{E}^{\times}$, where $f_{E_{v}}$ is a smooth function on $E_{v}^{\times}$(by Lemma 2) with $f_{E_{v}}(a \gamma)=\chi_{v}(a) f_{E_{v}}(\gamma)$ $\left(a \in F_{v}^{\times}, \gamma \in E_{v}^{\times}\right)$, which is spherical ( $R_{E_{v}}^{\times}$-invariant) for $v \notin V$, and for almost all $v$ it is the unit element: $f_{E_{v}}^{0}(\gamma)=\chi_{v}\left(\boldsymbol{\pi}_{v}\right)^{v a l_{E_{v}}(\gamma)}$ (if $v$ is non-split), and $f_{E_{v}}^{0}((a, b)$ ) equals 1 if $|a|_{v}=|b|_{v}$, and zero otherwise (if $v$ splits). The sum over $\mu$ ranges over all characters of $\mathbb{A}_{E}^{\times} / E^{\times}$whose restriction to $\mathbb{A}^{\times} / F^{\times}$is $\chi$, and

$$
\mu\left(f_{E}\right)=\int_{\mathbb{A}_{E}^{\times} / \mathbb{A}^{\times}} \mu(a) f_{E}(a) d a
$$

The measure da is such that $\int_{\mathbb{A}_{E}^{\times} / \mathbb{A}^{\times} E^{\times}} d a=2$, the Tamagawa number of $\operatorname{Res}_{E / F} \mathbb{G}_{m} / \mathbb{G}_{m}$.

Note that the sums over $\gamma$ and $\mu$ are equal by the Poisson summation formula. Since $f_{E}(\bar{a})=f_{E}(a)$, we have $\bar{\mu}\left(f_{E}\right)=\mu\left(f_{E}\right)$, where $\bar{\mu}(a)=\mu(\bar{a})$.

Lemmas 3 and 4 assert that at $v \notin V$, we have $\mu_{v}\left(f_{E_{v}}\right)=\operatorname{tr} I\left(\mu_{1 v}, \mu_{1 v}^{-1} ; f_{v}\right)$, where $\mu_{1 v}$ is related to $\mu_{v}$ as in the Theorem. On the other hand, in (1), $\pi_{v}\left(f_{v}\right)$ acts as
zero on $\varphi \in \pi$ unless $\varphi$ is $K_{v}$-invariant on the right, in which case $\pi_{v}\left(f_{v}\right)$ acts as multiplication by the scalar $\operatorname{tr} \pi_{v}\left(f_{v}\right)$. A standard argument of "generalized linear independence of characters" (see, e.g., [F2], p. 758), using the absolute convergence of our sums, simple unitarity estimates, and the Stone-Weierstrass theorem, implies the following. Put $\mathbb{K}(V)=\prod_{v \notin V} K_{v}$, and let $\pi^{\mathbb{K}(V)}$ be the space of $\mathbb{K}(V)$-invariant vectors in the space of $\pi$.
6. Proposition. Fix an unramified $G_{v}$-module $\pi_{v}^{*}$ for each $v \notin V$. For any $f_{v} \in$ $C_{c}^{\infty}\left(G_{v}\right), v \in V$, put $f=\left(\bigotimes_{v \in V} f_{v}\right) \otimes\left(\bigotimes_{v \notin V} f_{v}^{0}\right)$. Then $(1)+\sum_{\eta^{2}=\chi} I_{\eta}=\sum_{\mu} \mu\left(f_{E}\right)$, where in (1) the first sum ranges over the cuspidal representations $\pi$ of $\operatorname{PGL}(2, \mathbb{A})$ with $\pi_{v} \simeq \pi_{v}^{*}$ for all $v \notin V$, and the second sum is over a smooth orthonormal basis $\{\varphi\}$ for the spaces $\pi^{\mathbb{K}(V)}$. The sum over $\eta, \eta^{2}=\chi$, ranges over those characters $\eta$ with $I\left(\eta_{v}, 1 / \eta_{v}\right) \simeq \pi_{v}^{*}$ for all $v \notin V$. The sum over $\mu$ ranges over those characters of $\mathbb{A}_{E}^{\times} / E^{\times}$such that for $v \notin V$ the component $\mu_{v}$ is unramified, and defines the representation $I\left(\mu_{1 v}, \mu_{1 v}^{-1}\right)$, which is required to be equivalent to $\pi_{v}^{*}$.

By the Chebotarev density theorem the sum over $\mu$ consists of at most one pair $\{\mu, \bar{\mu}\}$ of non-zero contributions. Since every smooth function on $E_{v}^{\times}$which transforms under $F_{v}^{\times}$according to $\chi_{v}$ and whose values at $\gamma \in E_{v}^{\times}$and $\bar{\gamma}$ defer by a multiple of $\chi_{v}(-1)$, is obtained as $f_{E_{v}}$ from some $f_{v}$, for some $u_{v}$ and $w_{v}$, we conclude, on choosing $\pi_{v}^{*}=I\left(\mu_{1 v}, \mu_{1 v}^{-1}\right)(v \notin V)$, that for each $\mu$ as in the Theorem there exists (a unique) $\pi(\mu)$, as in the Theorem; it is the unique $\pi$ which occurs in (1), unless $\chi=\eta^{2}$ and $\mu=\bar{\mu}$, since the sum $\sum_{\mu} \mu\left(f_{E}\right)$ of Proposition 6 is non-zero. This $\pi=\pi(\mu)$ has the property that $R^{\chi}(\varphi) \neq 0$ for some $\varphi \in \pi$.

On the other hand, by the rigidity theorem for $G L(2)$ (see [JS]), at most one $\pi$ can contribute to the sum (1) of Proposition 6. Let $\pi$ be a cuspidal representation of $P G L(2, \mathbb{A})$ such that $\int_{\mathbb{Z}^{2} G \backslash G} \varphi_{1}(g) \Theta_{u}(g) \bar{\Theta}_{w}^{\chi}(g) d g$ is non-zero for some $u$ and $w$, and $\chi$, and a smooth form $\varphi_{1}$ in the space of $\pi$. We can choose a sufficiently large finite set $V$, and $\pi_{v}^{*}=\pi_{v}$ for $v \notin V$, such that the equality $(1)=\sum_{\mu} \mu\left(f_{E}\right)$ of Proposition 6 holds. The $I_{\eta}$ vanish again by [JS]. We may assume that the orthonormal basis of $\pi^{\mathbb{K}(V)}$ in (1) contains $\varphi_{1}$. Since $\pi$ is cuspidal, it is generic, namely there exists a form $\varphi_{2}$ in its space such that $W_{\varphi_{2}, \psi}(e) \neq 0$. We may assume that either $\varphi_{2}$ is $\varphi_{1}$, or $\varphi_{2}$ is orthogonal to $\varphi_{1}$. In any case, the space of endomorphisms of $\pi_{v}$ is spanned by the operators $\pi_{v}\left(f_{v}\right), f_{v} \in C_{c}^{\infty}\left(Z_{v} \backslash G_{v}\right)$. Hence we can choose $f_{v}(v \in V)$ such that $\prod_{v \in V} \pi_{v}\left(f_{v}\right)$ maps $\varphi_{2}$ to $\varphi_{1}$, and any vector in $\pi^{\mathbb{K}(V)}$ which is orthogonal to $\varphi_{2}$, to 0 . With this choice of $f$ in Proposition 6, the two sums of (1) consist of one term each. Our $\pi$, and $\varphi_{2}$, index the only possibly non-zero term:

$$
\int_{\mathbb{Z}^{2} G \backslash \mathbb{G}}\left(\pi(f) \varphi_{2}\right)(g) \Theta_{u}(g) \bar{\Theta}_{w}^{\chi}(g) d g \cdot \bar{W}_{\varphi_{2}, \psi}(e),
$$

which is non-zero by our choice of $\varphi_{2}$ and $\pi(f) \varphi_{2}=\varphi_{1}$. Since the sum (1) is nonzero, there is $\mu$ such that $\mu\left(f_{E}\right) \neq 0$, by the equality of Proposition 6 , and if $\chi$ is a square, $\mu$ satisfies $\mu \neq \bar{\mu}$. This proves that $\pi$ with $R^{\chi}(\varphi) \neq 0$ for some $\varphi \in \pi$ is necessarily of the form $\pi(\mu)$, and the Theorem follows.

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